



Crustal layering in northeastern Tibet: A case study based on joint inversion of receiver functions and surface wave dispersion

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Complete List of Authors:	Deng, Yangfan; Guangzhou Institute of Geochemistry, Chinese Academy of Sciences, State Key Laboratory of Isotope Geochemistry Shen, Weisen; University of Colorado at Boulder, Department of Physics Xu, Tao; Institute of Geology and Geophysics, Chinese Academy of Sciences, State Key Laboratory of Lithospheric Evolution Ritzwoller, Michael; University of Colorado at Boulder, Department of Physics
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4 1 **Crustal layering in northeastern Tibet: A case study based on joint**
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6 2 **inversion of receiver functions and surface wave dispersion**
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9 3 Yangfan Deng^{1,2}, Weisen Shen⁴, Tao Xu^{2,3}, and Michael H. Ritzwoller⁴

10 4 1 - State Key Laboratory of Isotope Geochemistry, Guangzhou Institute of Geochemistry, Chinese
11 Academy of Sciences, Guangzhou, 510640, China. yangfandeng@gig.ac.cn

12 5
13 6 2 - State Key Laboratory of Lithospheric Evolution, Institute of Geology and Geophysics,
14 Chinese Academy of Sciences, Beijing, 100029, China

15 7
16 8 3 - Chinese Academy of Sciences, Center for Excellence in Tibetan Plateau Earth Sciences,
17 Beijing, 100101, China

18 9
19 10 4 - Department of Physics, University of Colorado at Boulder, Boulder, CO 80309, USA.
20 weisen.shen@colorado.edu
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24 12 **Abstract:**

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26 13 Recently constructed models of crustal structure across Tibet based on surface wave data
27 14 display a prominent mid-crustal low velocity zone but are vertically smooth in the crust.
28 15 Using six months of broadband seismic data recorded at 22 stations arrayed
29 16 approximately linearly over a 440 km observation profile across northeastern Tibet (from
30 17 the Songpan-Ganzi block, through the Qaidam block, into the Qilian block), we perform
31 18 a Bayesian Monte Carlo joint inversion of receiver function data with surface wave
32 19 dispersion to address whether crustal layering is needed to fit both data sets
33 20 simultaneously. On some intervals a vertically smooth crust is consistent with both data
34 21 sets, but across most of the observation profile two types of layering are required: a
35 22 discrete low velocity zone (LVZ) or high velocity zone (HVZ) formed by two
36 23 discontinuities in the middle crust and a doublet Moho formed by two discontinuities
37 24 from 45-50 km to 60-65 km depth connected by a linear velocity gradient in the
38 25 lowermost crust. The final model possesses (1) a mid-crustal LVZ that extends from the
39 26 Songpan-Ganzi block through the Kunlun suture into the Qaidam block consistent with
40 27 partial melt and ductile flow and (2) a mid-crustal HVZ bracketing the South Qilian
41 28 suture coincident with ultrahigh pressure metamorphic rocks at the surface. (3)
42 29 Additionally, the model possesses a doublet Moho extending from the Qaidam to the
43 30 Qilian blocks which probably reflects increased mafic content with depth in the
44 31 lowermost crust perhaps caused by a vertical gradient of eclogitization. (4) Crustal
45 32 thickness is consistent with a step-Moho that jumps discontinuously by 6 km from 63.8
46 33 km (± 1.8 km) south of 35° to 57.8 km (± 1.4 km) north of this point coincident with the
47 34 northern terminus of the mid-crustal LVZ. These results are presented as a guide to future
48 35 joint inversions across a much larger region of Tibet.

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57 36 **Keywords:** Joint inversion, receiver functions, surface waves, Tibet, crust
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1. Introduction

The expansion of seismic instrumentation in Tibet has led to the rapid emergence of velocity models of the Tibetan crust and upper mantle. The emplacement of broadband seismometers, in particular, allows for the observation of surface waves based both on ambient noise (e.g., Yao et al., 2006; Yang et al., 2010; Zheng et al., 2010; Zhou et al., 2012; Karplus et al., 2013) and earthquake data (e.g., Caldwell et al., 2009; Feng et al., 2011; Li et al., 2013; Zhang et al., 2014). Studies based on surface waves provide information primarily about shear wave speeds in the crust and uppermost mantle beneath Tibet (e.g., Yao et al., 2008; Li et al., 2009; Duret et al., 2010; Huang et al., 2010; Guo et al., 2012; Yang et al., 2012; Li et al., 2013; Xie et al., 2013; Chen et al., 2014). A positive attribute of surface wave studies is that information is spread homogeneously across much of Tibet, but at the price of relatively low resolution both laterally and vertically. The vertical resolution of models derived from surface waves presents a particular challenge, as surface waves do not image discontinuities in seismic velocities well. Receiver functions image internal interfaces better than surface waves and there have been several studies based on them across Tibet (e.g., Zhu and Helmberger, 1998; Vergne et al., 2002; Wittlinger et al., 2004; Xu et al., 2007; Shi et al., 2009; Zhao et al., 2011; Sun et al., 2012; Yue et al., 2012; Tian and Zhang, 2013; Xu et al., 2013b; Tian et al., 2014; Zhang et al., 2014). Receiver functions, however, only provide information near seismic stations and less powerfully constrain structures between interfaces than surface waves (e.g., Ammon et al., 1990). The joint interpretation of receiver functions along

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4 58 with surface wave dispersion, however, provides information about vertical layering that
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7 59 surface waves alone may miss (e.g., Ozalabey et al., 1997; Julia et al., 2000; Bodin et al,
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10 60 2012; Shen et al. 2013a). Using data from the USArray in the US, Shen et al. (2013a)
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12 61 present a method to invert receiver functions and surface wave dispersion jointly based
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15 62 on a Bayesian Monte Carlo method to produce a model of shear wave speeds (and other
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18 63 variables) along with uncertainties in the crust and uppermost mantle. Shen et al.
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20 64 (2013b,c) show that vertically smooth crustal models can fit both data sets acceptably
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23 65 except in several regions and, therefore, across most of the western and central US the
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26 66 introduction of layering within the crystalline crust is not required to fit the receiver
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29 67 function data used in their study. The purpose of the current paper is to address the same
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31 68 questions for Tibet with a particular focus on northeastern Tibet: (1) Can surface wave
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34 69 dispersion and receiver function data be fit simultaneously with vertically smooth models
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37 70 in the crystalline crust? (2) If not, then what is the nature of the discontinuities between
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40 71 the sediments and Moho that must be introduced to allow both data sets to be fit
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43 72 simultaneously? (3) Finally, what do the answers to these questions imply about the
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46 73 thickness and structure of the crust in northeastern Tibet?
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48 74 In this paper we use six months of data (late 2010 to mid-2011) from a linear array of 22
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51 75 broadband seismometers deployed in the Songpan-Ganzi block, the Qaidam block, and
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54 76 the Qilian block of northeastern Tibet (**Fig. 1**, blue diamonds). We use these data to
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57 77 produce receiver functions (with uncertainties) at 23 evenly spaced geographical
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4 78 locations spanning a distance of 440 km along the “observation profile”. Using the
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7 79 method of Shen et al. (2013a), we jointly invert these receiver functions along with
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10 80 Rayleigh wave phase speed data taken from the study of Xie et al. (2013) to produce
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12 81 shear velocity models (with uncertainties) of the crust and uppermost mantle beneath the
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15 82 observation profile. We produce two models, one that varies smoothly vertically in the
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18 83 crystalline crust (**Model 1**) and another one (**Model 2**) that allows for crustal
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20 84 discontinuities that are adapted to the receiver functions. In contrast to the US (Shen et al.,
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23 85 2013b,c), we present evidence here that across most of the observation profile crystalline
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26 86 crustal discontinuities and/or a doublet Moho are needed to fit the receiver functions.
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29 87 The crust in northeastern Tibet has already been the subject of studies based on seismic
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32 88 reflection and refraction profiling (see **Fig. 1**) as well as receiver functions (e.g., Vergne
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35 89 et al., 2002; Shi et al., 2009; Zhao et al., 2011; Yue et al., 2012; Tian and Zhang, 2013;
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38 90 Xu et al., 2013b; Tian et al., 2014). These studies have observed significant variations in
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41 91 crustal structure (e.g., Karplus et al., 2011; Mechie et al., 2012) and thickness, including
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44 92 in some cases stepwise thickening of Moho (e.g., Zhu and Helmberger, 1998; Vergne et
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47 93 al., 2002; Wittlinger et al., 2004; Jiang et al., 2006). Understanding such variations is
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50 94 critical to test conflicting hypotheses related to the formation and evolution of the Tibetan
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53 95 plateau (e.g., Molnar et al., 1993; Tapponnier et al., 2001). Northeastern Tibet also is the
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56 96 site of choice to study remote effects of the India-Asia collision (Metivier et al., 1998;
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59 97 Meyer et al., 1998; Chen et al., 1999; Pares et al., 2003). In addition, the region contains
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4 98 the Caledonian orogeny and petrological and isotopic data point to high pressure and
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7 99 ultrahigh pressure metamorphism (UH/UHP; Liu et al., 2003; Luo et al., 2012), which is
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10 100 inferred from the presence of eclogite and garnet peridotite as well as coesite-bearing
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12 101 gneiss in the north Qaidam (Song et al., 1996, 2006; Yang et al., 2002). Such
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15 102 metamorphism may have been caused by the burial and exhumation of the metamorphic
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18 103 rocks in the uppermost mantle along a Paleozoic subduction zone (e.g., Yin et al., 2007).
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21 104 We are, nevertheless, unaware of previous studies based on the joint inversion of surface
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24 105 wave data and receiver function across northeastern Tibet. Other researchers have
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27 106 performed such joint inversions elsewhere in Tibet, notably in southeast Tibet or
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30 107 southwest China (e.g., Li et al., 2008; Sun et al., 2014; Wang et al., 2014; Bao et al., 2015)
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32 108 and the Lhasa terrane (e.g., Xu et al., 2013a). Thus, the models produced in this study
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35 109 both may guide future joint inversions at large scales across Tibet and also provide new
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38 110 information about the structure and thickness of the crust in northeastern Tibet.
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41 111 In Section 2, we discuss the receiver function and surface wave phase speed data sets that
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44 112 we use in the joint inversion. The hypothesis test to determine if crystalline crustal
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47 113 layering is needed and its characteristics are presented in section 3, in which we contrast
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50 114 model characteristics and data fit with and without intra-crustal layering. In Section 4, we
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53 115 discuss the implications of the final crustal velocity model (Model 2) in terms of
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56 116 mid-crustal partial melt in the Songpan-Ganzi block and its potential intrusion into the
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59 117 Qaidam block, the coincidence of a mid-crustal high velocity zone with HP/UHP rocks in
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4 118 the Qaidam block, the nature and location of the Moho doublets and the step-Moho along
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7 119 the observation profile, and finally compare our estimates of crustal thickness with earlier
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10 120 studies in the region.

121 **2. Data and Methodology**

122 Between November 2010 and June 2011, a passive seismic experiment was carried out by
123 the Institute of Geology and Geophysics, Chinese Academy of Sciences, from the
124 Songpan-Ganzi block to the Qilian block (**Fig. 1**). Twenty-two broadband seismographs
125 (Reftek-72A data loggers and Guralp CMG3-ESP sensors with 50 Hz-30 s bandwidth;
126 represented by the blue diamonds in **Fig. 1**) were deployed at intervals of about 20 km.

127 The profile covers the northeastern margin of the Tibetan plateau. The
128 Northwest-Southeast trending Animaqing-Kunlun-Muztagh suture (Kunlun fault) and the
129 South Qilian suture divide the profile into three principal geological units: the
130 Songpan-Ganzi block, the Qaidam block, and the Qilian block.

131 **2.1 Receiver Functions**

132 **2.1.1 Sensor Orientation**

133 Misorientation of sensors will cause amplitude errors in receiver functions (Niu and Li,
134 2011). Before computing the receiver functions, we attempt to determine station
135 orientation using P-wave particle motions (e.g., Niu and Li, 2011). A misorientation less
136 than 7° in azimuth is not expected to affect receiver functions or surface wave
137 polarizations significantly (Niu and Li, 2011). We found that only one of the 22 stations

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4 138 exhibited a misorientation azimuth larger than 7° , and corrected the orientation for this
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7 139 station.

10 140 **2.1.2 Calculation of Receiver Functions**

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12 141 Receiver functions are determined by deconvolving the vertical component seismogram
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15 142 from the radial component, thereby isolating the receiver site effects from other
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18 143 information contained in the teleseismic P waveforms (e.g., Ammon, 1991). We write this
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21 144 schematically in the frequency (ω) domain as follows:

$$23 \quad 145 \quad RF(\omega) = \frac{R(\omega)}{V(\omega)} \quad (1)$$

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27 146 where $R(\omega)$ is the radial component at a particular station, $V(\omega)$ is the vertical component,
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30 147 and $RF(\omega)$ is the receiver function which is typically displayed after it is transformed
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33 148 back into the time domain to produce $RF(t)$. In practice, however, after rotating the
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36 149 observed North and East components to the Radial and Transverse directions, we
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39 150 calculate the receiver functions using a time-domain iterative deconvolution method
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42 151 (Ligorria and Ammon, 1999). During this process, we apply a low-pass Gaussian filter to
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45 152 produce receiver functions with a dominant period of about 1 sec, thereby reducing
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48 153 high-frequency noise (and signal). Prior to this calculation, we selected teleseismic
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51 154 P-waveforms from earthquakes with magnitudes $M_w \geq 5.5$ in the epicentral distance
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54 155 range from $30^\circ - 90^\circ$ (**Fig. 1**, inset). We make corrections to the receiver functions in both
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57 156 time and amplitude by normalizing to a reference slowness of 0.06 deg s^{-1} (Jones and
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60 157 Phinney, 1998). Those receiver functions that have P wave slownesses greater than

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4 158 0.1 deg s⁻¹ or smaller than 0.04 deg s⁻¹ are discarded before the normalization. The
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7 159 Vp/Vs ratio is set to 1.75 in both the crust and mantle. The reason for this choice and its
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10 160 effects are discussed later in the paper.

11 12 13 161 **2.1.3 Quality Control**

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15 162 Following Shen et al. (2013a), we perform a three-step quality control process. Step 1:
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18 163 We remove receiver functions whose product with the vertical component seismogram
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21 164 poorly approximates the radial component. Step 2: We remove receiver functions with
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24 165 unrealistic amplitudes at zero time (greater than 1 or smaller than 0.02). Step 3: We
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26 166 employ a method known as ‘harmonic stripping’ (Shen et al., 2013a) to remove receiver
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29 167 functions that do not vary smoothly in azimuth. If j denotes the earthquake index, an
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32 168 observed receiver function at a particular station derives from a P wave that propagates at
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35 169 azimuth θ_j and is denoted $RF(\theta_j, t)$. In this step, we fit a truncated harmonic function to all
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38 170 such observed receiver functions from different earthquakes (i.e., azimuths) for each
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41 171 station at each time t as follows:

$$42 \quad H(\theta, t) = A_0(t) + A_1(t) \sin[\theta + \alpha_1(t)] + A_2(t) \sin[2\theta + \alpha_2(t)] . \quad (2)$$

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44 173 Here, the time functions A_i ($i = 0, 1, 2$) are the amplitudes of the three harmonic
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47 174 components of the receiver functions and the angles α_i are the initial phases for the
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50 175 azimuthally dependent components. This harmonic analysis is designed to identify the
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53 176 azimuthally smooth structural effects. If a given observed receiver function for
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56 177 earthquake j , $RF(\theta_j, t)$, disagrees with the harmonic fit $H(\theta, t)$ when $\theta_j = \theta$, we reject that

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4 178 receiver function. What remains are observed receiver functions that vary smoothly in
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7 179 azimuth.
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10 180 **Figure 2** presents an example of the result of the quality control process for station
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12 181 DKL21 whose location is identified in **Figure 3**. The original receiver functions from 360
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15 182 earthquakes are presented in **Figure 2a** separated by azimuth. Most receiver functions are
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18 183 from earthquakes at azimuths between about 40° and 200°, which are from the northeast
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21 184 to the south of the study region. Substantial disagreement amongst the receiver functions
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24 185 is apparent in **Figure 2a**. After quality control Steps 1 and 2, the number of receiver
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27 186 functions reduces to 149 as shown in **Figure 2b**. The 111 azimuthally smooth receiver
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30 187 functions that emerge from the harmonic analysis of Step 3 are shown in **Figure 2c**. After
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32 188 the quality control process is complete, we retain a total of 1145 receiver functions for the
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34 189 22 stations along the profile.

37 190 **2.1.4 Receiver Function CMCP Stacks**

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40 191 Shen et al. (2013a,b) advocated for the use of the function $A_0(t)$ from equation (2) as
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43 192 representative of the azimuthally independent structure beneath the station. However,
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46 193 **Figure 2** shows that the distribution of earthquakes in our study produces receiver
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49 194 functions that lie primarily in the azimuthal range from 40° to 200°, so $A_0(t)$ may be
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52 195 biased by azimuthally dependent structure near the station. **Figure 3** further illustrates
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54 196 this point by presenting the locations of the Moho piercing points of P waves (or P to S
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57 197 conversion points) retained after quality control. Moho piercing points are computed by
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4 198 ray tracing from each earthquake through a model with P velocities from IASP91
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7 199 (Kennett and Engdahl, 1991) but with crustal thickness from Xu et al. (2014). The
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10 200 piercing points are predominantly to the east and southeast of the stations at which the
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12 201 receiver functions are observed and are characteristic of structures there rather than near
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15 202 the stations. For this reason, we use harmonic stripping only for quality control and not to
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17 203 produce the stacked receiver functions at the stations. Rather, we stack receiver functions
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20 204 along the observation profile at a set of 23 stacking locations lying at 20 km intervals
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22 205 (**Fig. 3**, black dots, numbered 1-23). We stacked (i.e., averaged) the receiver functions
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25 206 with Moho piercing points lying within 0.15° of each stacking location. We refer to this
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28 207 as the Common Moho Conversion Point (CMCP) stacking method, which is somewhat
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31 208 similar to the CCP (Common Conversion Point) stacking method (e.g., Dueker and
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34 209 Sheehan, 1998). The weights used in stacking are shown in the inset panel in **Figure 3**
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36 210 showing nine square sub-boxes with sides of 0.1° . In each sub-box, we average all
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39 211 receiver functions with equal weight producing what we call a sub-box receiver function.
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42 212 Then the sub-box receiver functions are stacked (averaged) according to the weights
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45 213 presented in the inset panel where the central sub-box lies on the stacking location. The
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47 214 stacking locations together form what we call the “observation profile”.

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50 215 **Figure 4b,c** presents the stacked receiver functions along the observation profile with
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53 216 locations identified by the location numbers 1 through 23. **Figure 4b** shows the stacked
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56 217 receiver function waveforms themselves and **Figure 4c** shows the same information but
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4 218 with amplitude-dependent color shading. The P and P-to-S converted phases from the
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7 219 Moho can be seen clearly along the profile. The delay time between the P and P-to-S
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10 220 converted phases from the Songpan-Ganzi block (SB) to the Qilian block (QL) varies
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12 221 from about 7 to 8 sec. This delay time reduces northward along the stacking profile and
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15 222 becomes more complicated, showing a double peak at most locations north of stacking
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18 223 location 12. Additional complexities in the receiver functions also appear in **Figure 4**,
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20 224 which are discussed later in the paper.

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23 225 We also estimate uncertainties for each receiver function along the profile. First, we
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26 226 compute the standard deviation at each time among the receiver functions in each
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29 227 stacking sub-box. We then take the weighted average of these standard deviations to
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32 228 compute the uncertainty of the stacked receiver function, using the weights given in the
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35 229 inset panel of **Figure 3**. An example of these one standard deviation uncertainties can be
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37 230 seen for Location 13 as the grey shaded envelope in **Figure 5a**.

38 39 40 231 **2.2 Rayleigh wave phase velocity**

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42 232 Xie et al. (2013) mapped phase velocities (with uncertainties) across eastern Tibet and
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45 233 surrounding regions for Rayleigh (8–65 s) and Love (8–44 s) waves using ambient noise
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48 234 tomography based on data from the Program for Array Seismic Studies of the Continental
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51 235 Lithosphere (PASSCAL) and the Chinese Earthquake Array (CEArray). We interpolate
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54 236 the Rayleigh wave phase speed and uncertainty curves beneath the stacking points, an
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56 237 example of which is shown in **Figure 5b** for Location 13. Rayleigh wave phase speeds
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4 238 increase from about 3.1 km/s at 8 sec period to about 3.85 km/s at 65 sec period, and the
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7 239 uncertainty also increases with the period. Other example Rayleigh wave phase speed
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10 240 curves are presented later in the paper.

11 12 13 241 **2.3 Joint inversion of Receiver Functions and Surface Wave Dispersion**

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15 242 Internal interfaces such as sedimentary basement and Moho are difficult to resolve based
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18 243 on the inversion of surface wave data alone. While surface wave dispersion constrains
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21 244 well the vertically averaged velocity profile, it only weakly constrains velocity interfaces
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23 245 and strong velocity gradients. Receiver functions have complementary strengths to
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26 246 surface wave data (e.g. Ozalaybey et al., 1997; Julia et al., 2000; Du et al, 2002; Li et al.,
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29 247 2008) and the joint inversion of surface wave dispersion with receiver functions may be
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31 248 more reliable than structures derived exclusively on either data set alone (e.g., Julia et al.,
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34 249 2003, 2005; Chang and Baag, 2005; Shen et al., 2013a,b). Shen et al. (2013a) developed a
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37 250 non-linear Bayesian Monte-Carlo algorithm to estimate a Vs model by jointly
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40 251 interpreting Rayleigh wave dispersion and receiver functions as well as associated
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43 252 uncertainties. We apply this method here. We apply stacked receiver functions in the 0–11
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46 253 sec time band and Rayleigh wave phase speeds between 8 and 65 sec period at 20 km
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49 254 intervals along the observation profile. This time band and period range provides
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52 255 information about the top 80 km of the crust and uppermost mantle. In the models
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55 256 presented here, both the inversion of surface wave data performed by Xie et al. (2013)
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58 257 and the joint inversions of surface wave data and receiver functions presented for the first
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4 258 time here, we apply a V_p/V_s ratio of 1.75 in both the crust and uppermost mantle. We
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7 259 choose to fix this ratio primarily for consistency with the starting model (Xie et al., 2013).
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10 260 There is no doubt, however, that V_p/V_s varies with depth and along our observation
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12 261 profile. The V_p/V_s ratio trades off with crustal thickness and structures within the crust
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15 262 and, therefore, the depth to Moho and the amplitude and depth of structural features in
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18 263 the crust will depend on this choice.

20 21 264 **3. Crustal Structure Along the Profile**

22 23 24 265 **3.1 Smooth Starting Model from Surface Waves Alone**

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26 266 We start with the model of Xie et al. (2013), which is determined from Rayleigh and
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29 267 Love wave phase speed measurements alone determined from ambient noise tomography.
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32 268 (For the background to this study see: Shapiro et al., 2004; Bensen et al., 2007; Lin et al.,
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35 269 2008, 2009; Zheng et al., 2008; Yang et al., 2010, 2012; Ritzwoller et al., 2011; Zheng et
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38 270 al., 2011; Zhou et al., 2012). The model is composed of three layers stacked vertically
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41 271 with variable thicknesses but the crystalline crust is vertically smooth. The top layer is
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43 272 the sediments, which are isotropic ($V_s = V_{sv} = V_{sh}$) with constant velocity vertically.
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46 273 The middle layer is the crystalline crust, which is radially anisotropic ($V_{sv} \neq V_{sh}$). Each
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49 274 of V_{sv} and V_{sh} is given by five B-splines in the crystalline crust. The bottom layer is the
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52 275 radially anisotropic uppermost mantle in which V_{sv} is given by five B-splines and the
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55 276 difference between V_{sh} and V_{sv} is taken from an earlier model of the region (Shapiro
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58 277 and Ritzwoller, 2002; Shapiro et al., 2004). Sedimentary thickness and Moho depth were
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4 278 free variables in the inversion for this model, which applied several constraints including
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7 279 vertical crustal smoothness and positive jumps at the base of the sediments and crust. For
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10 280 the purposes here, we only use the V_{sv} part of the model at all depths and set the model
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12 281 to be isotropic ($V_s = V_{sv} = V_{sh}$) because we only invert Rayleigh waves and receiver
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15 282 functions. The V_p/V_s ratio both in the crust and mantle is set to 1.75.

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18 283 A plot of V_{sv} as a function of depth along the observing profile is presented in **Figure 6a**.
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21 284 In this model the crust thins slowly and continuously to the north from about 61.5 km in
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23 285 the Songpan-Ganzi block to 51.5 km in the Qilian block. More prominently, in the
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26 286 Songpan-Ganzi block mid-crustal V_{sv} is very slow, much slower than in the Qilian block.
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29 287 Hacker et al. (2014) argues that such slow shear velocities must be caused by partial melt
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31 288 in the middle crust.

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34 289 Although receiver functions were not used in the construction of the model of Xie et al.,
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37 290 we compute synthetic receiver functions and present them in **Figure 7a,b**, designed to be
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40 291 compared with the observed receiver functions in **Figure 4b,c**. The timing of the Moho
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42 292 P-to-S converted phases on the synthetic receiver functions is similar to the observed
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45 293 receiver functions but other aspects of the synthetics and observations are quite different.
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48 294 First, the positive swing on the synthetic P-to-S converted phase is too broad, which is
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51 295 caused by a strong vertical velocity gradient both above and below the Moho in the
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53 296 model. Second, internal crustal structures are reflected in the observed receiver functions
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55
56 297 that are entirely missing in the synthetics. Such structures are apparent in **Figure 7c**,

298 which presents the difference between the observed and synthetic receiver functions.

299 Third, there are also complexities in the observed receiver functions near the P-to-S
 300 converted phase north of 35° latitude that are not apparent in the synthetics. We measure
 301 reduced χ^2 misfit on the interval between t_i and t_f for each of the 23 stacking locations as
 302 follows:

$$C_{location}^2 = \frac{1}{t_f - t_i} \int_{t_i}^{t_f} \frac{(RF^{obs}(t) - RF^{pred}(t))^2}{S^2(t)} dt \quad (3)$$

304 where RF^{obs} and RF^{pred} are the observed and predicted receiver functions at the location,
 305 respectively, σ is the standard deviation at the location, and we take $t_i = 2$ sec and $t_f = 8$
 306 sec. These location specific reduced χ^2 values are then averaged over the 23 locations to
 307 determine the total reduced χ^2 , which is 5.1 for the starting model. These results suggest,
 308 not surprisingly, that there are complexities in the structure of the crust that are missing in
 309 the vertically smooth crustal model of Xie et al.

310 **3.2 Joint Inversion of Surface Waves and Receiver Functions with a** 311 **Vertically Smooth Crystalline Crust: Model 1**

312 To begin to model the complexities in crustal structure implied by the receiver functions
 313 and inferred in section 3.1, we first perform a joint inversion of the Rayleigh wave
 314 dispersion data and the receiver functions at each location along the profile but continue
 315 with the constraint that the model is vertically smooth in the crystalline crust. We refer to
 316 this model as having resulted from the vertically smooth joint inversion or as Model 1.

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4 317 Data such as those shown in **Figure 5a,b** are inverted jointing using the method
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7 318 described by Shen et al. (2013a,b). The starting model is the model of Xie et al. (2013)
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10 319 and we adopt the parameterization of this model with three modifications: (1) the model
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12 320 is isotropic at all depths (there is no radial anisotropy) such that $V_s = V_{sv}$, (2) we use
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15 321 seven B-splines for V_{sv} in the crust rather than five, and (3) we represent sedimentary
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18 322 velocities as a linear monotonically increasing function of depth rather than a constant.
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20 323 Importantly, as crystalline crustal structure is represented with B-splines, although larger
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23 324 in number than in the model of Xie et al., in this inversion the crystalline crust still is
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26 325 constrained to be smooth vertically. In section 3.3, this constraint is broken in order to
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29 326 introduce internal crustal interfaces that appear to be needed in order to fit the receiver
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31 327 function data.

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34 328 We perform the inversion with a Bayesian Monte Carlo method aimed to fit the Rayleigh
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37 329 wave dispersion and receiver functions jointly and equally well at each location along the
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40 330 observation profile. Uncertainty estimates in each type of data weight the relative
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43 331 influence of each data type in the likelihood function (i.e., misfit) and the inversion
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45
46 332 results in a posterior distribution of models that fit the data acceptably at each depth.
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48 333 **Figure 5c** shows the results of the inversion at Point 13 along the observation profile,
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50 334 presented with a gray corridor that represents the full width of the posterior distribution at
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52
53 335 each depth. The blue line is the mean of the posterior distribution at each depth.
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56 336 The mean value of the posterior distribution for V_{sv} at each depth along the observation
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4 337 profile is presented in **Figure 6b**. Compared with the starting model from Xie et al. (2013)
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7 338 determined from surface wave data alone in **Figure 6a**, vertical variations in Model from
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10 339 1 are sharper; e.g., the mid-crustal velocities in the Songpan-Ganzi block are confined to
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12 340 a narrower depth range, are slower, and are overlain by a thin veneer of higher velocities
13
14
15 341 at about 10 km depth, there are higher velocities in the lowermost crust (50-60 km)
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17
18 342 bracketing the Kunlun fault, and the mid-crustal velocities in the Qilian block are
19
20 343 generally faster although very low velocities appear in the lowermost crust south of the
21
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23 344 South Qilian suture.
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26 345 **Figure 8a,b** presents the synthetic receiver functions computed from Model 1. A
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29 346 comparison, in particular, between the observed receiver functions in **Figure 4c** and the
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32 347 synthetics in **Figure 8b** illustrates the improvement in fit via the introduction of vertically
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35 348 smooth internal crustal structures that nevertheless produce receiver function arrivals
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38 349 between the direct P arrival and the P-to-S conversion. **Figure 8c** quantifies this
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41 350 comparison by presenting the difference between the observed and synthetic receiver
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44 351 functions. Contrasting **Figure 8c** with **Figure 7c** shows that the fit to the receiver
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47 352 functions is greatly improved compared with the model of Xie et al. even though the
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50 353 model remains vertically smooth in the crust. The overall χ^2 , defined by equation (3), is
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53 354 1.4, which represents a 72% reduction in the variance relative to the starting model.
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56 355 Therefore, the introduction of receiver functions in the joint inversion is advisable even
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59 356 when retaining a vertically smooth model.
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4 357 Nevertheless, there remain considerable differences between the observed and synthetic
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6 358 receiver functions, particularly in the boxes marked (A) and (B) in **Figure 8c**. This can be
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9 359 seen more clearly in **Figure 9**, which presents vertical profiles beneath Point 6 (box A)
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11
12 360 and Point 18 (box B) from Model 1 as the red lines in **Figure 9a** and **9d**, respectively.
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15 361 (Blue lines and the corridor of accepted models are discussed later, in section 3.3.)
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17 362 **Figure 9a,d** (red lines) illustrate how the receiver functions at these two points are misfit
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20 363 by Model 1. The misfit in the receiver function near Point 6 is somewhat subtle, but at
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22
23 364 Point 18 the double peak between 6 and 8 seconds cannot be fit with this model and
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26 365 neither can the swings in the receiver function between 3 and 5 seconds.
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28 366 For these reasons, it is necessary to move beyond vertically smooth crystalline crustal
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31 367 models in order to fit the receiver function data in Tibet at least in some (perhaps most)
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34 368 locations. Interfaces within the crust are needed, therefore, to fit the receiver function
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37 369 data in detail. This is a different conclusion than drawn by Shen et al. (2013b,c) for the
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40 370 western and central US, where vertically smooth crystalline crustal models were found to
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43 371 suffice to fit surface wave dispersion and receiver function data jointly except in isolated
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46 372 areas across this region. Of course, the crust is much thinner in the US than in our region
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48
49 373 of study.

374 **3.3 Joint Inversion of Surface Waves and Receiver Functions with a** 375 **Layerized Crystalline Crust: Model 2**

376 To move beyond the vertically smooth crystalline crustal model from the joint inversion

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4 377 presented in section 3.2 (**Figs. 6b, 8**), we introduce different mid-crustal discontinuities in
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7 378 Regions 1, 2, and 3, which are identified in **Table 1**.

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10 379 Region 1 encompasses latitudes between about 33.6° and 34.2° , locations numbered 5-7,
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12 380 which lie in the northern part of the Songpan-Ganzi block to the Kunlun fault. In this
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15 381 region we introduce two mid-crustal discontinuities to the starting model, one at 20 km
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18 382 depth and one at 40 km depth and allow a constant velocity perturbation between them.
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21 383 The depths of these discontinuities and the amplitude of the perturbation are introduced
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23 384 as free variables in the inversion. The result at Point 6, which is contained in Region 1, is
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26 385 shown in **Figure 9c**. The grey envelope denotes the full width of the posterior distribution,
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28
29 386 the blue line marks the mean of the posterior distribution at each depth, and the red line is
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31
32 387 the mean of the posterior distribution of Model 1 (section 3.2). The introduction of these
33
34 388 three degrees of freedom acts to restrict the depth extent of the low velocity zone (LVZ)
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37 389 in the central crust, increase the shear wave speed across most of the lower crust, and
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39
40 390 reduce crustal thickness relative to Model 1. The result is a considerably better fit to the
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42 391 receiver function (blue line in **Fig. 9a**), particularly the P-to-S Moho conversion phase
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44
45 392 that appears near 8 seconds.

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47 393 Region 2 lies between latitudes of about 35.6° and 36.2° , locations numbered 17-20,
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49
50 394 which is the northern part of the Kunlun block to the southern Qilian block,
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53 395 encompassing the Southern Qilian suture. In this region the receiver functions are more
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56 396 complicated than elsewhere along the profile, and we introduce six degrees of freedom to
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4 397 the starting model. First, we introduce two mid-crustal discontinuities at 30 km and 40
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7 398 km depth and allow a constant velocity perturbation between them. These three degrees
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10 399 of freedom allow for a high velocity zone (HVZ) in the central crust to develop. Second,
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12 400 we also allow for a “doublet Moho” by introducing three more degrees of freedom to
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14 401 produce a linear velocity gradient in the lowermost crust (or uppermost mantle) with
15
16
17 402 variable depth and upper and lower shear velocities. The result at Point 18, which is
18
19
20 403 contained in Region 2, is shown in **Figure 9f**. A HVZ is introduced in the middle crust
21
22 404 between depths of about 30 to 40 km and there are two prominent discontinuities that
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25 405 compose the doublet Moho, one near 45 km and another nearer to 60 km depth with a
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28 406 linear velocity gradient between these depths. The result is a much better fit to the
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31 407 receiver function (blue line in **Fig. 9d**), including the double P-to-S Moho conversion
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34 408 phase that appears between 6 and 8 seconds, the positive swing near 3.5 seconds, and the
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37 409 negative swings near 2.5 and 4.5 seconds.

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39 410 Finally, Region 3 comprises a discontinuous set of locations numbered 1-2, 14-16, and
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41 411 21-22. In these locations we allow for a doublet Moho. Two of these ranges of points
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44 412 bracket Region 2, which also contains a doublet Moho, and the third occurs at the
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46
47 413 southern end of the observation profile. The locations with a doublet Moho are made
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49
50 414 clearer later in **Figure 11**, discussed later in the paper.

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53 415 The receiver functions in **Figure 10a,b** computed with the introduced mid-crustal
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56 416 discontinuities of Model 2 fit the observed receiver functions in **Figure 4b,c** much better

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4 417 than either the starting model or Model 1, particularly in Box B. The total reduced χ^2 (eqn.
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7 418 (3)) is 0.9, which is a 82% variance reduction relative to the starting model and a 36%
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10 419 variance reduction relative to Model 1. The residual is small across most of the profile
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12 420 with the principal exception at times greater than about 7 seconds near the south Qilian
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15 421 suture (Box B), which we believe may be due to further layering in the uppermost mantle
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17
18 422 near the Qilian block.

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21 423 Model 2, the model from the joint inversion with the layers introduced in the crystalline
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23 424 crust, is shown in **Figure 6c**. The low velocity zone near 34° latitude in the
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25
26 425 Songpan-Ganzi block has been accentuated further by lowering the minimum velocities
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28
29 426 and the uppermost and lowermost crust has correspondingly been made faster. More
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31 427 substantially, the model between 35.5° and 36.2° latitude, bracketing the South Qilian
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34 428 suture, now has a high velocity zone introduced near a depth of 35 km with a doublet
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37 429 Moho (as shown).

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40 430 In summary, the surface wave dispersion data and receiver functions can be fit with at
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42 431 smooth crustal model across part of the observation profile, principally between location
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44
45 432 numbers 8-14 in the middle of the observation profile, but not in Regions 1-3 (**Table 1**).
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48 433 In these regions, crustal discontinuities must be introduced to fit the receiver functions. In
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51 434 Region 1, this produces a vertically narrower LVZ with a lower shear wave speed
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53 435 minimum. Region 2 is more complicated, requiring a HVZ in the middle crust. Beneath
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56 436 Regions 2 and 3 a doublet Moho at depths of about 50 km and 60 km provides a
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4 437 significant improvement in fit to the receiver functions.
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7 438 **4. Discussion**
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10 439 **4.1 Mid-Crustal LVZ in the Songpan-Ganzi Block: Evidence for Partial**
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12 440 **Melt**
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15 441 Crustal low velocity zones have been identified across Tibet by a number of studies (e.g.,
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17 442 Kind et al., 1996; Cotte et al., 1999; Rapine et al., 2003; Shapiro et al., 2004; Xu et al.,
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19 443 2007; Caldwell et al., 2009; Guo et al., 2009; Li et al., 2009; Yao et al., 2008, 2010; Acton
20
21 444 et al., 2010; Jiang et al., 2011). Yang et al. (2012) summarizes evidence from surface
22
23 445 waves for a mid-crustal low velocity zone (LVZ) across much of Tibet. Such evidence
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25 446 generally supports the internal deformation model of Tibetan evolution where the
26
27 447 medium is treated as a non-rigid continuum (e.g., England and Houseman, 1986; England
28
29 448 and Molnar, 1997) and may particularly favor ductile “channel” flow in the middle and/or
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31 449 lower crust (e.g., Bird, 1991; Clark and Royden, 2000; Searle et al., 2011). Based on the
32
33 450 more recent model of crustal shear velocities of Xie et al. (2013), our starting model
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35 451 presented along the observation profile in **Figure 6a**, Hacker et al. (2014) argue that the
36
37 452 low mid-crustal shear velocities across Tibet are indicative of partial melt. In the region
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39 453 of study, if we identify the LVZ as shear wave speeds (V_{sv}) below 3.4 km/s in the middle
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41 454 crust, then the LVZ extends from the Sonpan-Ganzi block through the Kunlun fault into
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43 455 the Qaidam block as far north as 34.9° (**Fig. 6a**). In this region, the V_p/V_s ratio was
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45 456 identified by Xu et al. (2014) to be greater than 1.75. As elsewhere in Tibet, block or
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4 457 terrane boundaries do not appear to obstruct crustal LVZs in the middle to lower crust
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7 458 (e.g., Yang et al., 2012; Jiang et al., 2014). Here, we ask the question whether the
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10 459 introduction of receiver functions in Region 1 (identified in **Table 1**) in the inversion
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12 460 increases or decreases the likelihood of partial melt in the middle crust.
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15 461 First, in the results of the joint inversion of Rayleigh wave dispersion and receiver
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18 462 functions with an imposed crustal vertical smoothness constraint (Model 1, **Fig. 6b**), the
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21 463 shear velocities in the middle crust actually rise compared to the starting model
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23 464 constructed with surface wave data alone (**Fig. 6a**). This somewhat reduces the likelihood
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26 465 of partial melt in the middle crust. This rise occurs because the attempt to fit the receiver
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29 466 functions with a vertically smooth crystalline crust increases crustal thickness, which
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31
32 467 reduces predicted Rayleigh wave speeds in the period band sensitive to the middle and
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34 468 lower crust and the increased mid-crustal shear wave speeds compensate. Minimum V_{sv}
35
36
37 469 speeds in the middle crust beneath the Songpan-Ganzi block in Model 1 mostly lie
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39 470 between 3.3 and 3.4 km/s but are somewhat lower in the starting model (**Fig. 6a**).
41
42 471 However, when two mid-crustal discontinuities are introduced in the joint inversion
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44
45 472 (Model 2, **Fig. 6c**) in order to improve the fit to the receiver function data, then the LVZ
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47
48 473 is accentuated in the middle crust beneath the Songpan-Ganzi block. In particular, the
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50 474 transitions to the LVZ from above and below are sharper, the LVZ is confined to a
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53 475 narrower depth range (20-40 km as opposed to 15-45 km), and the minimum shear wave
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56 476 speeds are lowered by about 0.1 km/s, which makes partial melt more somewhat more
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4 477 likely than in the model of Xie et al. (2013).
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7 478 Consequently, improving the fit to receiver functions by introducing internal crustal
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10 479 discontinuities modifies the shape and nature of LVZ beneath northern Tibet but does not
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12 480 reduce the likelihood of mid-crustal partial melt. Rather, it results in a slight increase in
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14
15 481 the likelihood of mid-crustal partial melt.
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18 482 **4.2 Mid-Crustal HVZ in the Qilian Block: Coincident with HP/UHP**

19 483 **Metamorphism**

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24 484 The Qilian Caledonian orogenic belt is believed to be the product of the convergence and
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26 485 collision between the North China Craton with the Qilian and Qaidam terranes during the
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28
29 486 Early Paleozoic Era (Yang et al., 2001). Ultrahigh Pressure (UHP) metamorphic rock are
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32 487 found along and near the South Qilian suture and several geological and tectonic models
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35 488 have been proposed to explain the origin of these rocks (Wang and Chen, 1987; Yang et
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37 489 al., 1994, 1999, 2000a,b, 2002; Song et al., 2006, 2009; Yin et al., 2007; Zhang et al.,
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39
40 490 2009). The crust near the South Qilian suture is known to be geophysically complicated,
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42 491 possessing highly variable V_p/V_s ratios (e.g., Xu et al., 2014), high and variable residual
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45 492 Bouguer gravity anomalies (EGM2008, Pavlis et al., 2012), and complicated receiver
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48 493 functions (e.g., Vergne et al., 2002; Xu et al., 2014). Our crustal model adds to this
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51 494 picture of crustal complexity by introducing a prominent high velocity anomaly at a
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53 495 depth of about 35 km that brackets the South Qilian suture (latitudes from 35.6° to 36.2°)
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56 496 directly below and adjacent to surface outcrops of UHP metamorphic rocks. We believe
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4 497 that the most likely interpretation is that this anomaly results from compositional
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7 498 heterogeneity, presumably of relatively enriched mafic rocks. Whether and how this
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10 499 anomalous structure relates to the UHP metamorphic rocks of the areas remains an area
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12 500 for further investigation.

15 501 **4.3 The “Doublet Moho”: Evidence for a Transitional Lower Crust**

18 502 **Bracketing the South Qilian Suture**

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21 503 A doublet Moho has been observed in earlier studies in at least two different locations
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24 504 beneath the Lhasa Terrane in southern Tibet (Kind et al., 2002; Nabelek et al., 2009; Li et
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26 505 al., 2011) and has been interpreted by Nabelek et al. (2009) to be caused by eclogitized
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29 506 lower crust from the Indian Plate underplating the Tibetan crust. As **Figure 6c** shows, we
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32 507 infer a doublet Moho to bracket the South Qilian Suture at latitudes from about 35.1° to
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34
35 508 36.5°. Part of the doublet Moho underlies the mid-crustal high velocity body discussed in
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38 509 section 4.2. The depth to both discontinuities that compose the doublet Moho are
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40 510 presented more clearly in **Figure 11**.

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43 511 The doublet Moho extends from depths of between 45-50 km to 55-65 km and
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46 512 encompasses an anomalously strong vertical velocity gradient. We interpret the latter
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49 513 discontinuity as classic Moho because beneath it lie shear wave speeds consistent with
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52 514 mantle rocks. Within the transition zone between these two discontinuities, shear wave
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55 515 speeds lie between about 3.8 km/s and 4.2 km/s. The high end of this range is only
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57 516 slightly faster than the lower crust south of this region where there is a single Moho, in
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4 517 the Qaidam and Songpan-Ganzi blocks; thus, the higher velocities rise up to shallower
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7 518 depths beneath the doublet Moho than further south. Thus, we do not find evidence that
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10 519 the lower crust encompassed by the doublet Moho is compositionally distinct from the
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12 520 lower crust elsewhere along the observation profile, but the lower crustal composition
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15 521 extends to shallower depths.
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18 522 The cause of the doublet Moho is not clear. One possibility is that there is no distinct
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21 523 crust – mantle division, but rather crustal and mantle rocks are interlayered in this region.
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23 524 Searle et al. (2011) proposes that the principal mineralogical composition of the Tibetan
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26 525 lower crust is granulite and eclogite with some ultramafic restites. Yang et al. (2012)
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29 526 argue that shear wave speeds of eclogite are expected to be about 4.4 km/s at the
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32 527 temperature and pressure conditions of the lower crust within Tibet. In situ lower crustal
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35 528 shear wave speeds all along the observing profile are significantly lower than this value
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38 529 so that the lower crust may be only partially eclogitized if eclogitization does indeed
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41 530 occur. Schulte-Pelkum (2005) estimates that 30% of the lower crust undergoes
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44 531 eclogitization in southern Tibet. Thus, an increasing fraction of eclogite versus granulite
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47 532 with depth may explain the vertical velocity gradient in the lower crust beneath the entire
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50 533 observation profile. What is unusual and requires further study is that the high shear wave
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53 534 velocities are compressed into a much narrower depth range in the doublet Moho regions
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56 535 than elsewhere along the observation profile.

536 **4.4 Crustal Thickness and Stepwise Crustal Thickening: Comparison**

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4 537 **with Previous Studies**
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7 538 The crust and upper mantle in different parts of the northeastern Tibetan plateau have
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10 539 been studied by a range of controlled source seismic experiments (e.g., Zhang et al.,
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12 540 2011b), many of which are identified in **Figure 1**. Our estimate of crustal thickness and
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15 541 its uncertainty are presented in **Figure 11** as red error bars and two lines. The dashed red
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18 542 line appears where we estimate a doublet Moho and indicates the shallower of the two
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21 543 discontinuities. The lower of the two discontinuities is indicated with a solid line and we
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24 544 take this to indicate crustal thickness. Crustal thickness averages 63.8 km (± 1.8 km) south
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26 545 of about 35° latitude and 57.8 km (± 1.4 km) north of this latitude, where the listed
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29 546 uncertainties are the standard deviation of the mean values south and north of this latitude.
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31 547 In fact, our results are consistent with a step-Moho at about 35° latitude. The location of
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34 548 the step is not coincident with the Kunlun fault, but is located about 50 km north of it.
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37 549 Rather it appears to be related to the termination of the mid-crustal LVZ that we find
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40 550 extends from the Songpan-Ganzi block into the Kunlun block as have other researchers
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42 551 (e.g., Jiang et al., 2014). Electromagnetic studies have also found that mid-crustal high
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45 552 conductivity features interpreted as melt extend north beyond the Kunlun fault (Le Paper
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48 553 et al., 2012).

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50 554 Other studies have also inferred discrete steps in Moho in northern Tibet based on
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53 555 receiver function studies; some are considerably west of our observation profile (e.g.,
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56 556 Zhu and Helmberger, 1998) but others are quite close (e.g. Vergne et al., 2002). Vergne et
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4 557 al. argue that the stairsteps in Moho are located beneath the main, reactivated Mesozoic
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7 558 sutures in the region and take this as evidence against partial melt in the middle crust. We
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10 559 find, however, that the Moho step lies between the main sutures within our observation
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12 560 profile and appears to coincide with a change in middle crustal structure. In fact, it
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15 561 appears to lie near the northern edge of the mid-crustal LVZ, which we follow Hacker et
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18 562 al. (2014) to interpret as being caused by partial melt in the middle crust. We posit,
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20 563 therefore, that the stairstep structure of Moho is consistent with a ductile middle crust and
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23 564 partial melt in the Tibetan crust.

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26 565 **Figure 11** presents crustal thickness estimates from other studies in the region for
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29 566 comparison with ours. Crustal thickness from the surface wave inversion of Xie et al.
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31 567 (2013), our smooth starting model which was based exclusively on the surface wave data
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34 568 we use here, slowly and continuously thins northward but is everywhere about five km
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37 569 thinner than our estimates as shown by the grey dotted line in **Figure 11**. The
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40 570 introduction of receiver functions causes the crust in our model to thicken along the entire
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43 571 observation profile relative to the starting model and bifurcate into a thicker southern
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46 572 zone that steps discontinuously to a thinner northern zone.

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48 573 Xu et al. (2014) used two methods to estimate crustal thickness based on the P wave data
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51 574 we use to produce receiver functions: PS migration and H-k stacking. We average their
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54 575 crustal thickness estimates and present them in **Figure 11** as the blue line. These
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57 576 estimates typically agree within one standard deviation with our results but vary more
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4 577 smoothly with latitude and do not as clearly show the step-Moho that we estimate. Two
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7 578 other cross-profiles, the MQ-JB (Liu et al., 2006) and the ALT-LMS (Wang et al., 2013)
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10 579 shown in **Figure 1**, exhibit similar crustal thickness at the intersections with our profile.
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12 580 There are greater differences with the active source crustal thickness estimates of Zhang
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15 581 et al. (2010), presented as the green line in **Figure 11**, which is more nearly constant with
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18 582 latitude.

20 583 **5. Conclusions**

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23 584 The results presented here highlight the significance of crustal layering in Tibet and the
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26 585 importance of parameterizing such layering in models of the Tibetan crust. Although on
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29 586 some intervals along the observation profile a vertically smooth crust is consistent with
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32 587 both data sets, across most of the observation profile two types of layering are required.
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34 588 First, there is the need for a discrete low velocity zone (LVZ) or high velocity zone (HVZ)
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37 589 formed by two discontinuities in the middle crust. Second, there is also the need for a
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40 590 doublet Moho formed by two discontinuities from 45-50 km to 60-65 km depth
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43 591 connected by a linear velocity gradient in the lowermost crust.

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45 592 After modifying the model parameterizing by introducing these structural variables, we
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48 593 find that the final model (Model 2) possesses the following characteristics. (1) The model
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51 594 has a mid-crustal low velocity zone that extends from the Gongpan-Ganzi block through
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54 595 the Kunlun suture into the Qaidam block consistent with partial melt and ductile flow. (2)
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56 596 There is also a mid-crustal high velocity zone bracketing the South Qilian suture that is
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4 597 coincident with ultrahigh pressure metamorphism of surface rocks that are believed to
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7 598 reflect deep crustal subduction in the Paleozoic. (3) Additionally, the model possesses a
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10 599 doublet Moho extending from the Qaidam to the Qilian blocks that probably reflects
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12 600 increased mafic content with depth in the lowermost crust perhaps caused by a gradient
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15 601 of eclogitization. (4) Crustal thickness is consistent with a step-Moho that jumps
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18 602 discontinuously from 63.8 km (± 1.8 km) south of 35° to 57.8 km (± 1.4 km) north of 35°,
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21 603 coinciding with the northern terminus of the mid-crustal LVZ that penetrates through the
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23 604 Kunlun suture into the Qaidam block.

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26 605 We present these results as a guide to future joint inversions across a much larger region
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29 606 of Tibet. As long as crustal models are suitably parameterized, historical data sets from
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32 607 PASSCAL and CEArray deployments, such as those employed by Yang et al. (2012) and
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35 608 Xie et al. (2013), as well as new deployments can be used for the joint inversion of
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38 609 surface wave data and receiver functions to reveal more accurate crustal structures across
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20 965 **Table 1.** Locations and types of crustal discontinuities.

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Region Number	Structures Introduced	Location Numbers	Latitude Range
1	Slow Mid-crustal Layer	5-7	33.6°-34.2°
2	Moho Doublet + Fast Mid-crustal Layer	17-20	35.6°-36.2°
3	Moho Doublet	1-2	33.1°-33.3°
		14-16	35.1°-35.6°
		21-22	36.2°-36.5°

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4 970 **Figure caption**
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7 971 **Figure 1.** The inset map presents locations of the distribution of teleseismic earthquakes
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9 972 used in this study. Blue triangles are the locations broadband seismometers used in this
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11 973 study; green stars are earlier seismic stations from the Lhasa-Golmud and Yushu-Gonghe
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13 974 profiles. Deep seismic sounding profiles and seismic reflection profiles include the
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15 975 following: HZ-JT: Hezuo-Jingtai profile (Zhang et al., 2013); MB-GD: Moba-Guide
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17 976 profile (Zhang et al., 2011); ALT-LMS: Altyn Tagh-Longmenshan profile (Wang et al.,
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19 977 2013); MK-GL: Markang-Gulang profile (Zhang et al., 2008); MQ-JB: Maqin-Jingbian
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21 978 profile (Liu et al., 2006); A: Galvé et al., 2002; B: Wang et al., 2011. Geological features
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23 979 include: ATF: Altyn Tagh fault; BNS: Bangong-Nujiang suture, JS: Jinsha suture, AKMS:
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25 980 Animaqing-Kunlun-Muztagh suture (or Kunlun fault), SQS: South Qilian suture, and the
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27 981 Songpan-Ganzi, Qaidam block, and Qilian blocks are identified. The region with
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29 982 Ultra-High Pressure (UHP) metamorphism is identified by the grey rectangle (Yang et al.,
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31 983 2002).
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44 985 **Figure 2.** Example of the quality control process for receiver functions (RFs) at station
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46 986 DKL21. (a) The full set of 360 observed RFs are plotted versus back azimuth. (b) The
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48 987 149 residual RFs after quality control steps 1 and 2. (c) The final 111 RFs after harmonic
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50 988 stripping, quality control step 3, in which only RFs that vary smoothly with azimuth are
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52 989 retained.
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4 991 **Figure 3.** Red triangles mark the locations of the 22 stations along the observation profile
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7 992 and green crosses show Moho piercing (or conversion) points (red crosses) of the
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10 993 incident P-waves. Blue triangles are Moho piercing points for station DKL21 (identified)
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12 994 The black dots indicate the stacking locations, which are separated by 20 km and
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15 995 numbered 1-23 (shown). The inset box contains the weights used to stack the receiver
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18 996 functions.

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23 998 **Figure 4.** (a) The elevation along the observation profile. (b) Moho conversion point
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25 999 (MCP) stacked receiver functions are illustrated with red waveforms as a function of the
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28 1000 stacking location number. (c) Smoothed color-coded image of the receiver functions.

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33 1002 **Figure 5.** Example of data and inversion result at location number 13. (a) The observed
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36 1003 receiver function (with uncertainty) is presented as the grey envelope. (b) The observed
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39 1004 Rayleigh wave phase speed curve (with uncertainties) is plotted with one standard
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42 1005 deviation error bars. (c) The full (grey) envelope of accepted models in the posterior
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45 1006 distribution from the joint inversion of the receiver function and dispersion data with a
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48 1007 smoothly varying crystalline crust (i.e., Model 1). The blue lines in (a) and (b) show the
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51 1008 predicted data from the best fitting model and the blue line in (c) presents the mean of the
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54 1009 posterior distribution at each depth.

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58 1011 **Figure 6.** The three models discussed here, all are V_{sv} (km/s). (a) The starting model
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4 1012 from Xie et al. (2013) constructed with Rayleigh wave dispersion data alone. (b) Model 1,
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7 1013 which results from the joint inversion of receiver functions and Rayleigh wave dispersion
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10 1014 without discontinuities in the crystalline crust. (c) Model 2, which results from the joint
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12 1015 inversion of receiver functions and Rayleigh wave dispersion with discontinuities in the
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15 1016 crystalline crust at the locations specified in **Table 1**.

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20 1018 **Figure 7.** (a) The computed receiver functions (RFs, red lines) from the starting model of
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23 1019 Xie et al. (2013). (b) A smoothed color-coded image of the computed RFs. (c) The
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26 1020 difference between computed RFs and observed RFs.

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30 1022 **Figure 8.** Similar **Fig. 7**, but here receiver functions are computing using Model 1, which
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33 1023 results from the joint inversion of receiver functions and Rayleigh wave dispersion
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36 1024 without discontinuities in the crystalline crust. The boxes denoted (A) and (B) identify
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39 1025 areas in which we particularly seek to improve the fit to the receiver functions.

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44 1027 **Figure 9.** Similar to **Fig. 5** in which example data and fits to receiver functions are
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47 1028 presented for two sample points: 6 (left column) and 18 (right column), which lie in the
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50 1029 boxes marked (A) and (B) in **Fig. 8**. Fits to the observed receiver functions and Rayleigh
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53 1030 wave phase velocities by both Model 1 and 2 are presented as red and blue lines,
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55 1031 respectively, in (a), (b), (d), and (e). Red and blue lines in (c) and (f) represent the best
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58 1032 fitting model of Model 1 (red) and Model 2 (blue). The full envelope of accepted models

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4 1033 in the inversion with crustal discontinuities (Model 2) is shown in (c) and (f).
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7 1034 **Figure 10.** Similar **Figs. 7** and **8**, but here RFs are computing using Model 2, which
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9 1035 results from the joint inversion of receiver functions and Rayleigh wave dispersion with
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11 1036 specified discontinuities in the crystalline crust. The boxes denoted (A) and (B) are
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13 1037 described in **Fig. 8**.
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20 1039 **Figure 11.** Our estimates of crustal thickness are presented with red error bars (1 standard
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22 1040 deviation). Where we infer a doublet Moho the lower interface is interpreted as Moho
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24 1041 (red solid line) and the upper interface is identified with the red dashed line. Crustal
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26 1042 thickness from Xie et al. (2013) is presented with grey dots, from the receiver function
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28 1043 study of Xu et al. (2013) is presented with the blue line, and from the deep seismic
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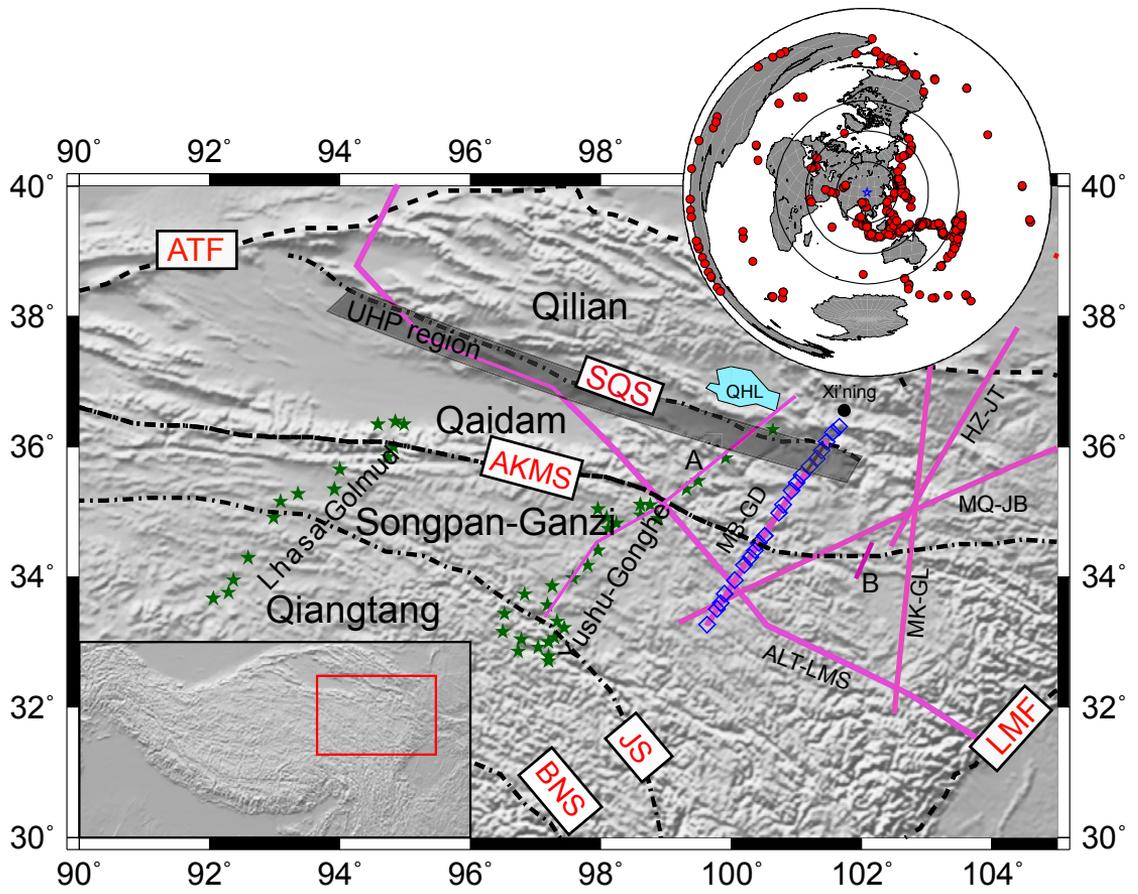
30 1044 sounding study of Zhang et al. (2011) with the green line. The symbols (diamond and
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32 1045 triangle) mark crustal thickness estimates crossing lines (Liu et al., 2006; Wang et al.,
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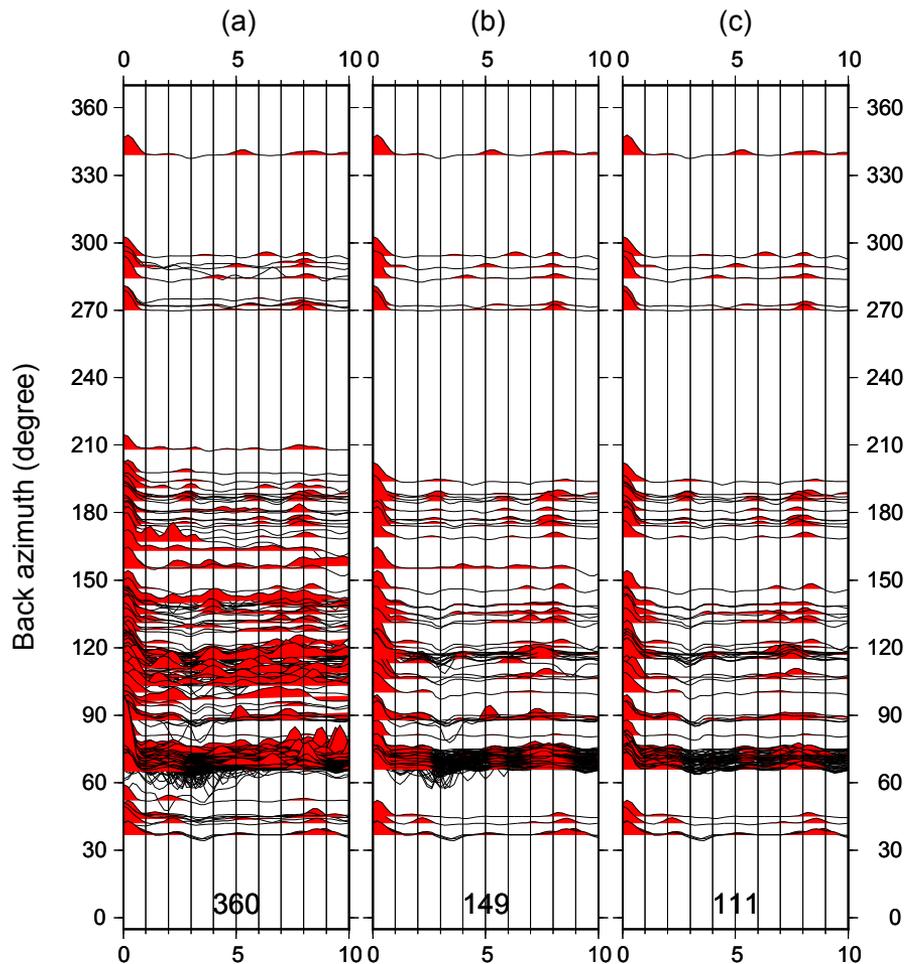
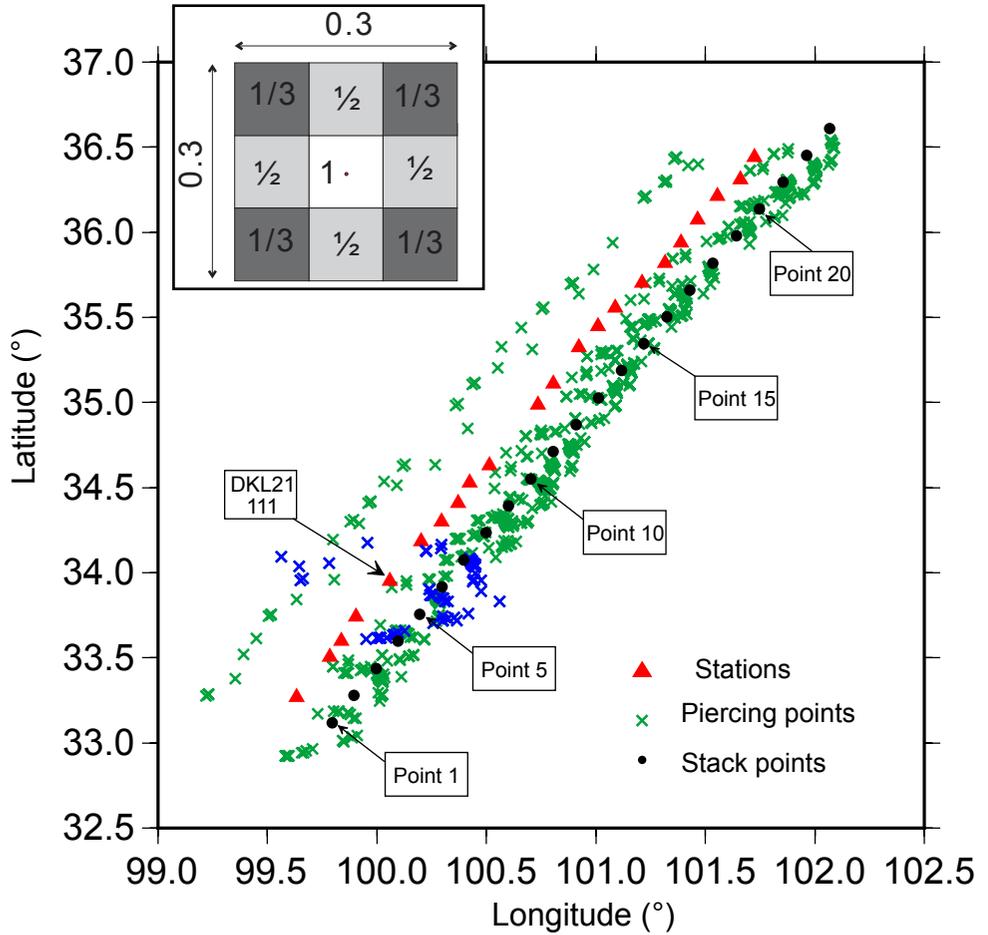
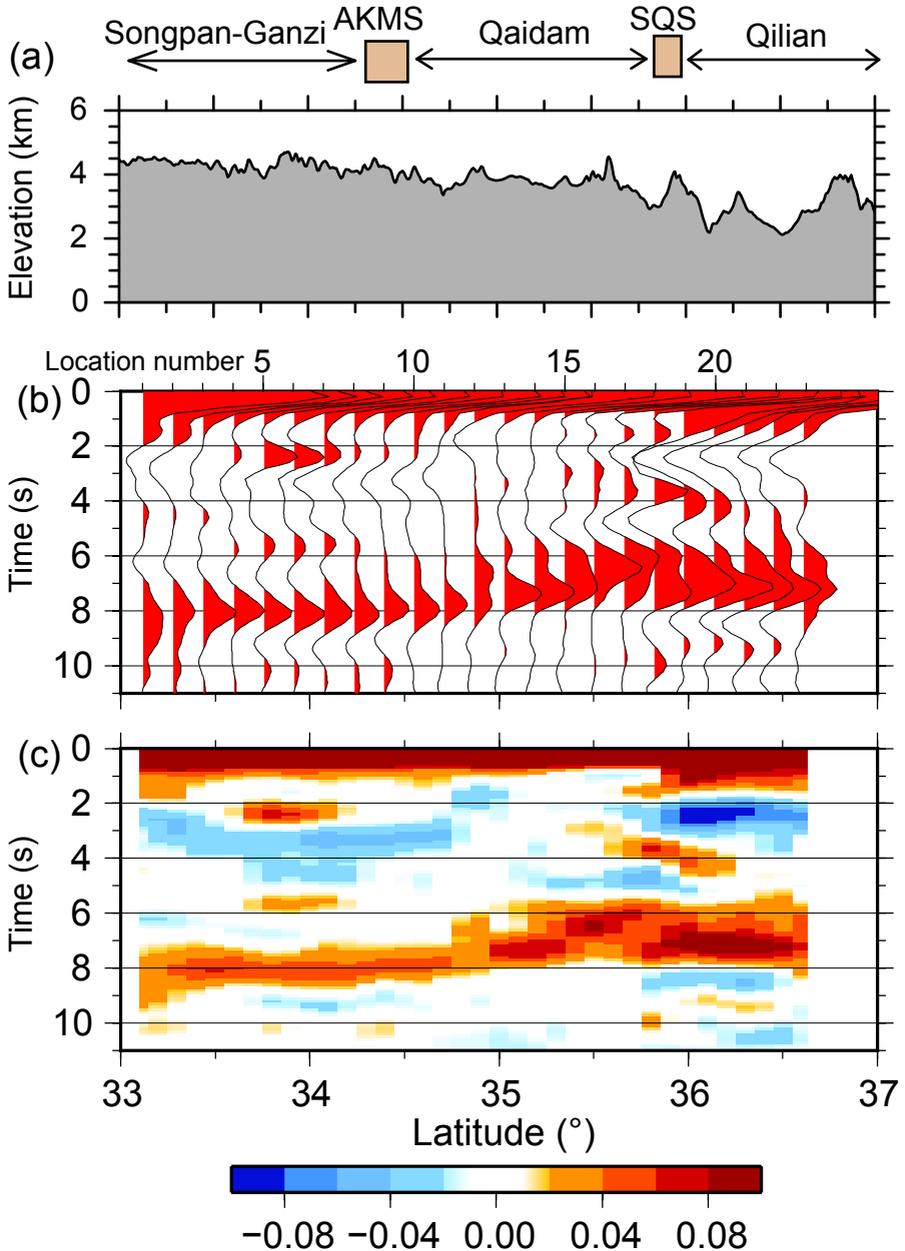


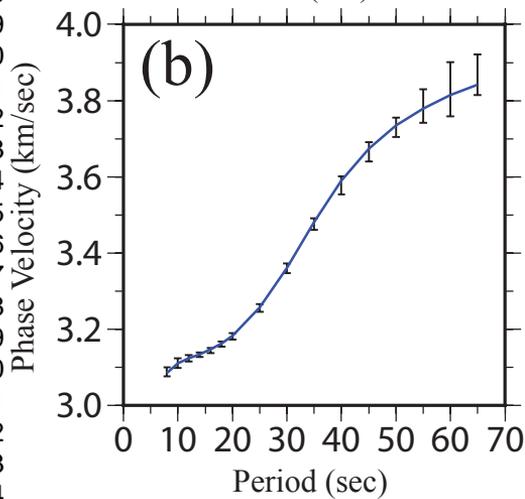
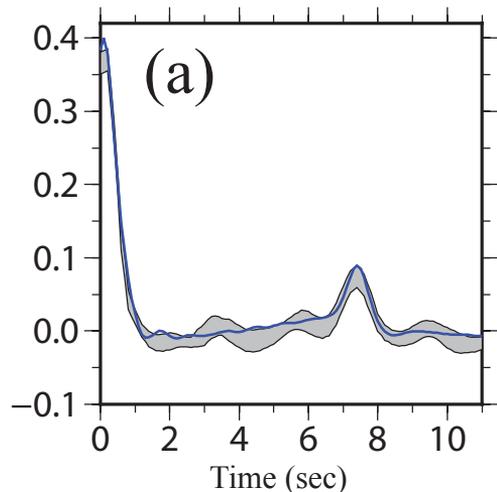
Figure 3



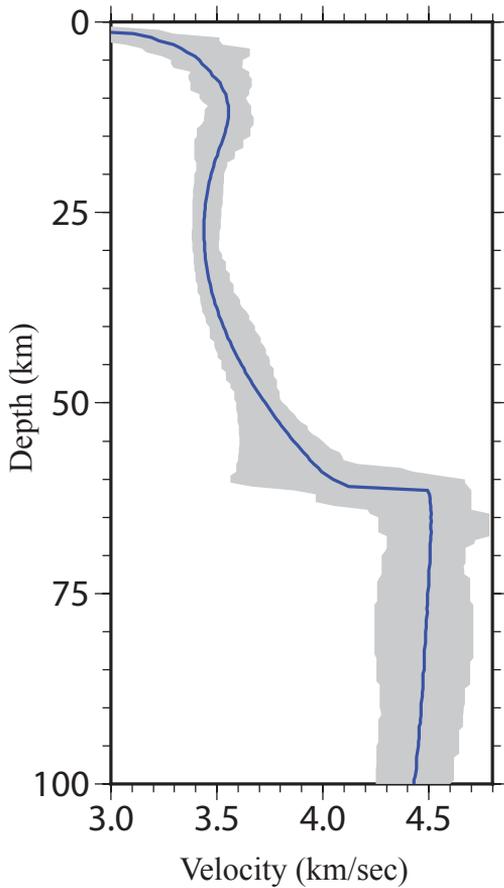


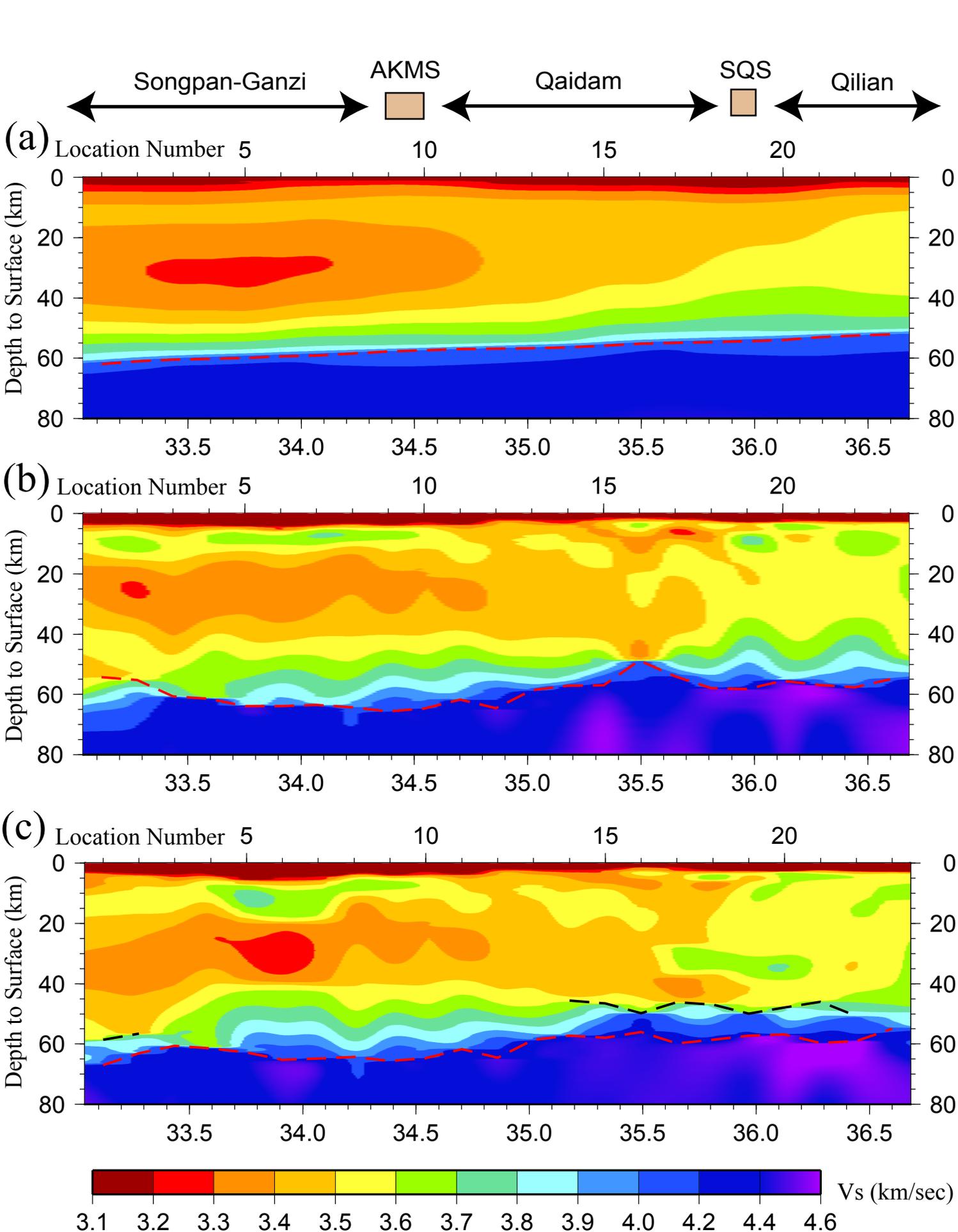
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Figure 5

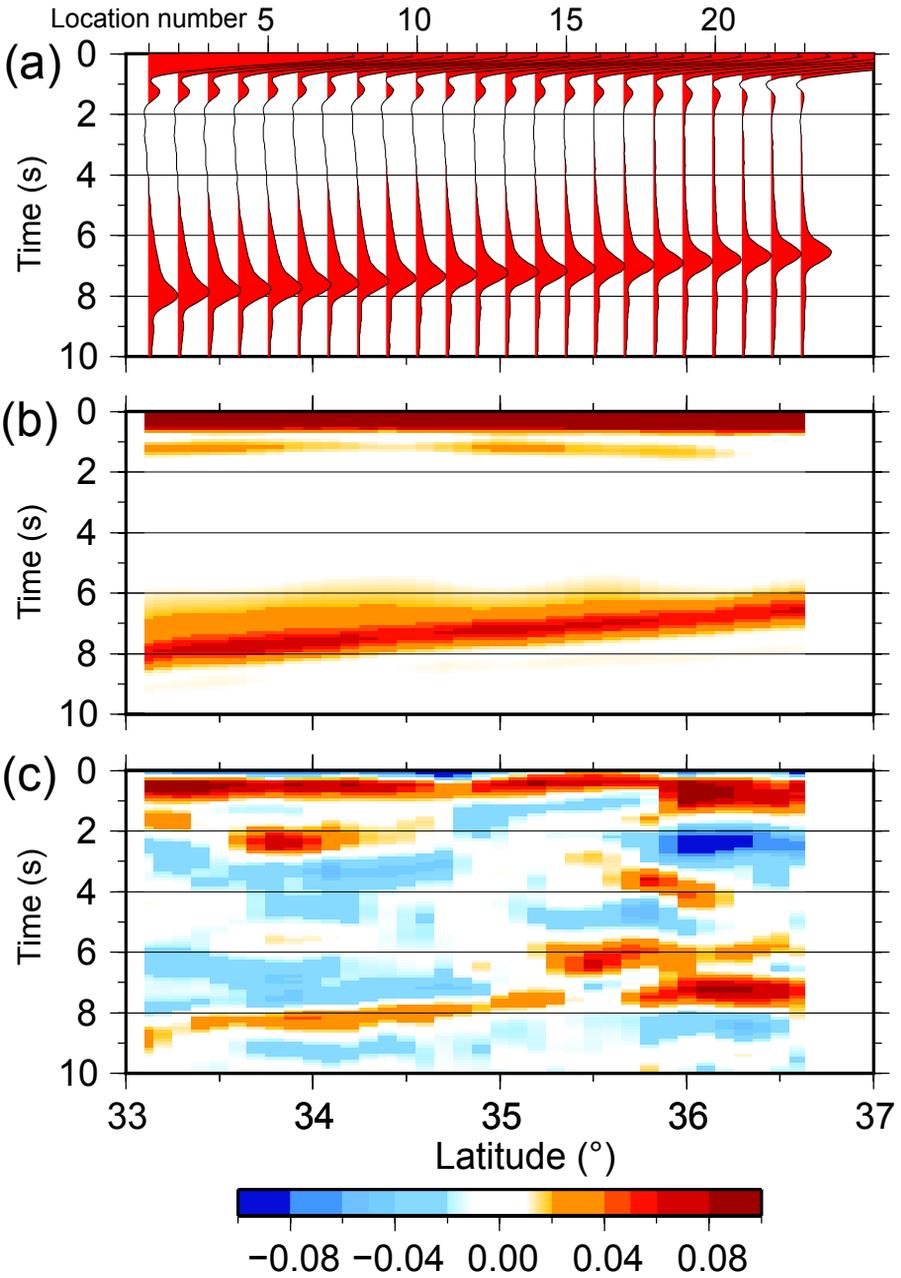
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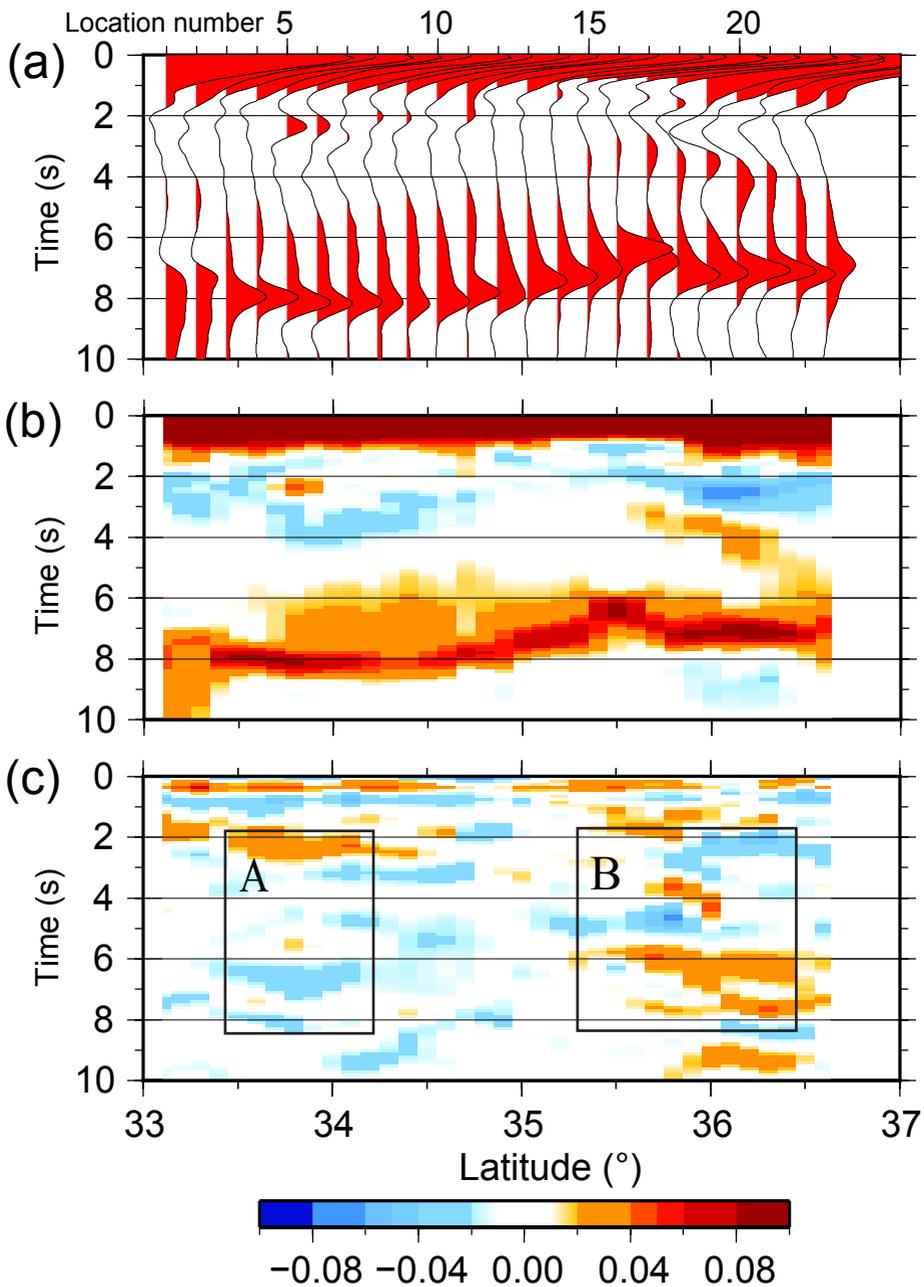
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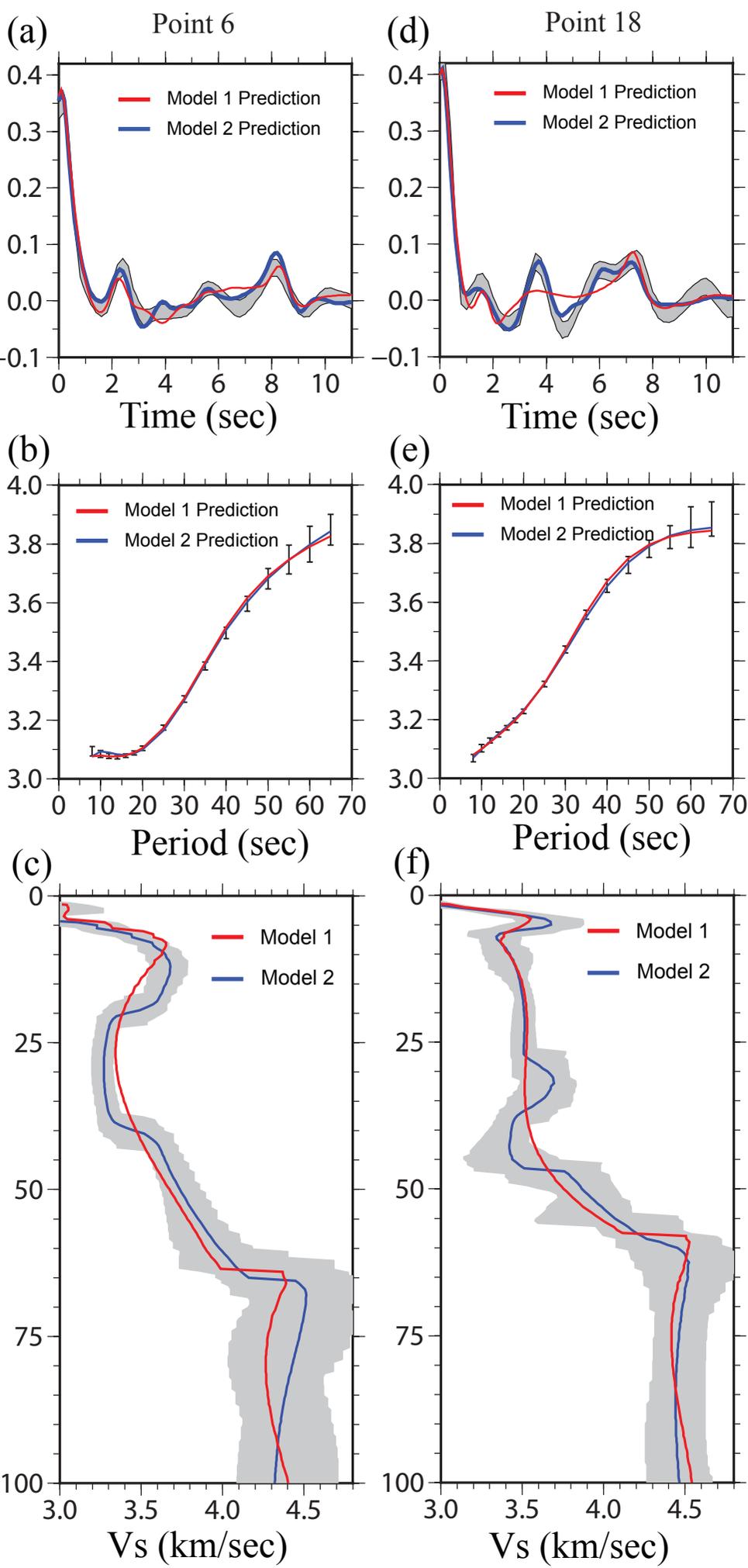
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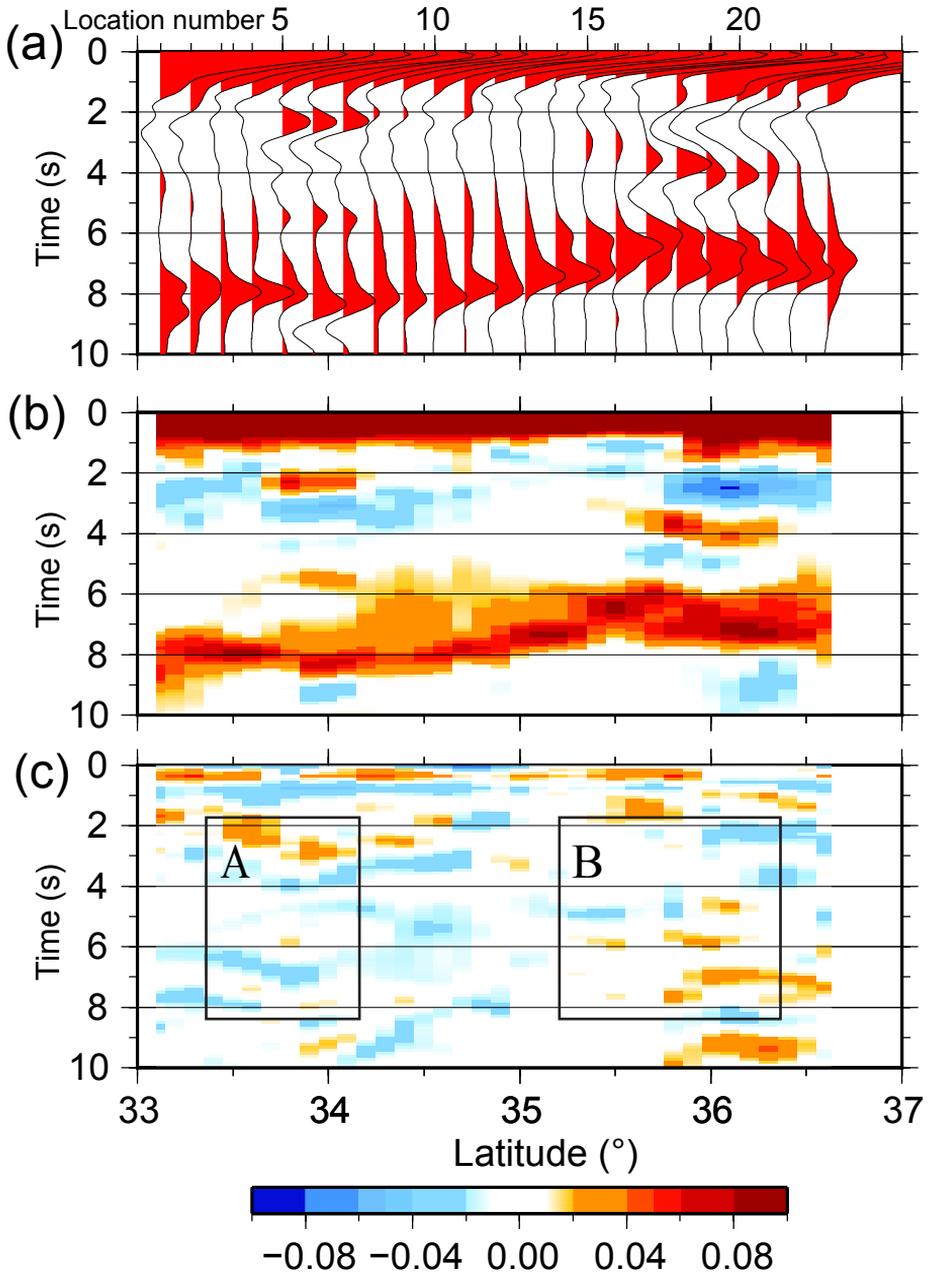




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